

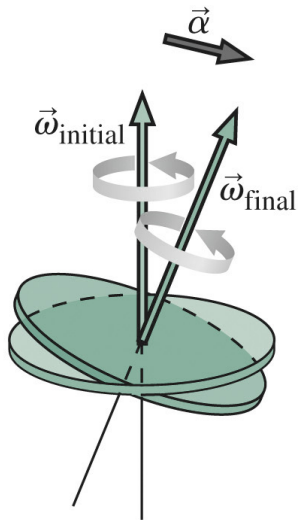
University Physics 226N/231N Old Dominion University More Rotational Motion

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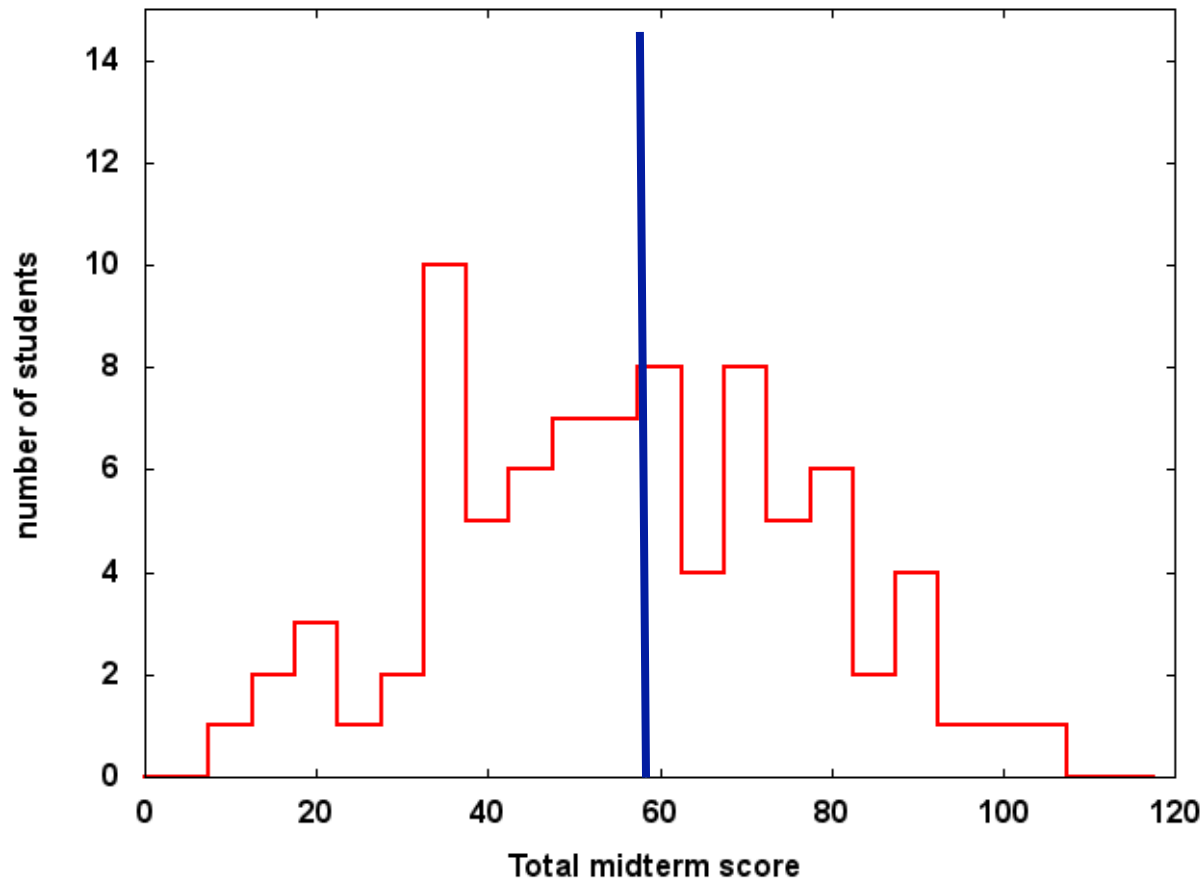
Friday October 26, 2012

Happy Birthday to Keith Urban, Jim Butcher, Natalie Merchant, and Pat Sajak!
Happy Frankenstein Friday and National Mincemeat Day!



Midterm 2 Grade Distribution

- Midterm 2 class average score: 58.2 (+/- 21.2 std deviation)



- Full statistics are in solution guide on class website
- Remember that I drop one midterm grade for final grading!



Rotational Inertia and the Analog of Newton's Law

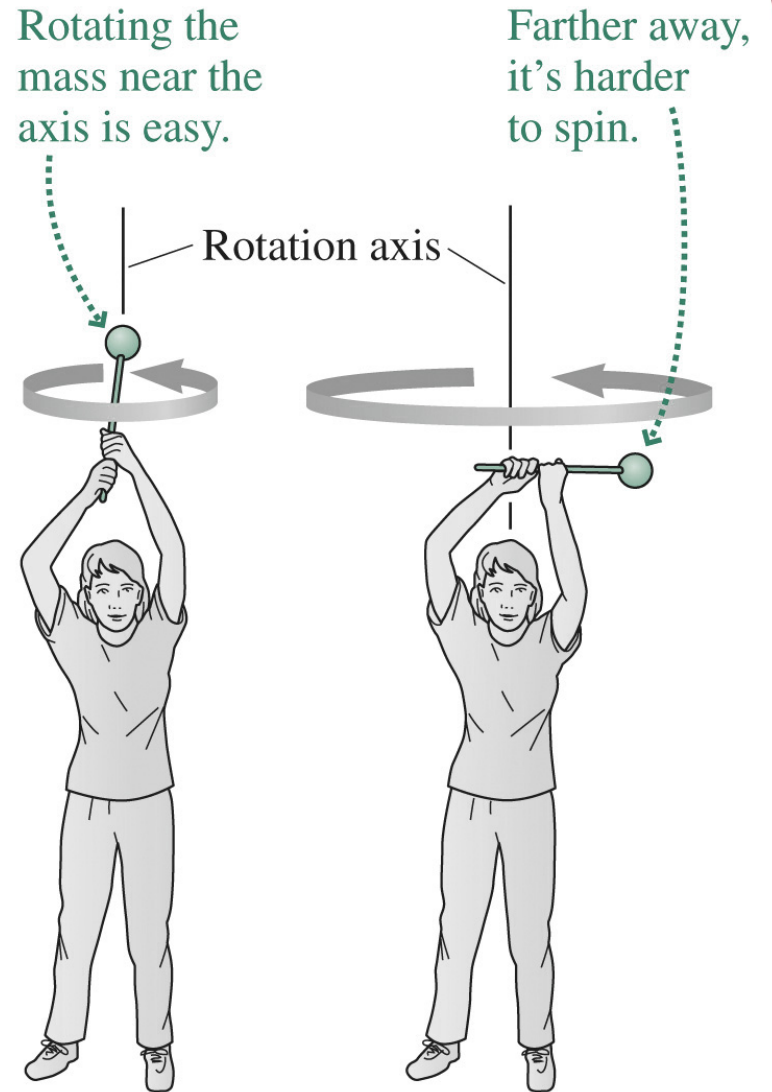
- **Rotational inertia** I (or moment of inertia) is the rotational analog of mass.
 - Rotational inertia depends on the distribution of mass and its distance from the rotation axis, similar to center of mass.
- Rotational acceleration, torque, and rotational inertia combine to give the rotational analog of Newton's second law $F = ma$

$$\tau = I\alpha$$

(or, more properly with vectors)

$$\vec{\tau} = I\vec{\alpha}$$

like $\vec{F} = m\vec{a}$



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Calculating Rotational Inertia

- For a single point mass m , rotational inertia is the product of mass with the square of the distance r from the rotation axis:

$$I = mr^2$$

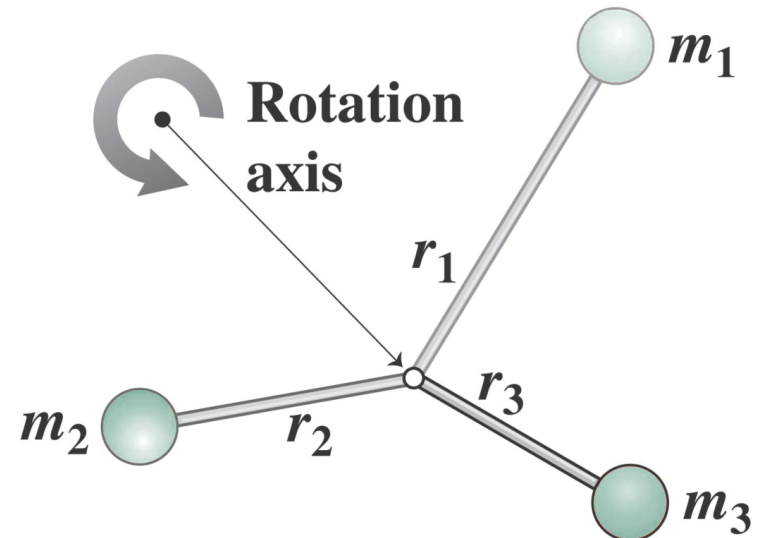
- For a system of discrete masses, the rotational inertia is the sum of the rotational inertias of the individual masses:

$$I = \sum m_i r_i^2$$

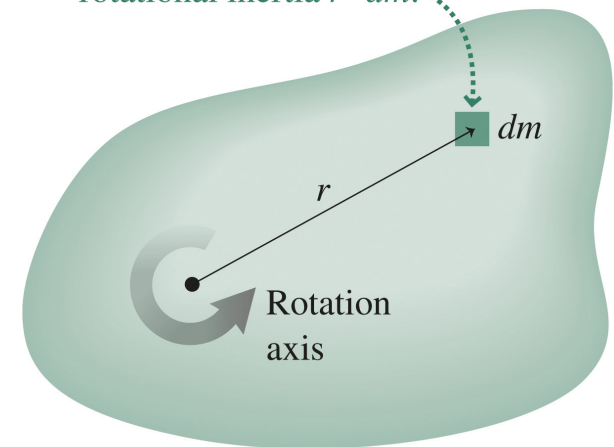
- For continuous matter, the rotational inertia is given by an integral over the distribution of matter:

$$I = \int r^2 dm$$

Similar to center of mass: $\vec{r}_{\text{cm}} = \frac{\int \vec{r} dm}{M}$

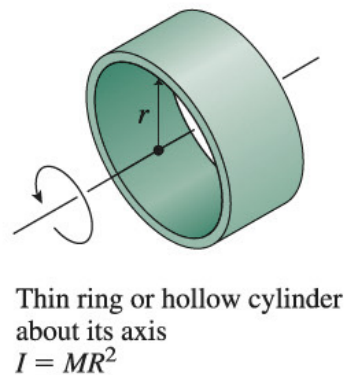
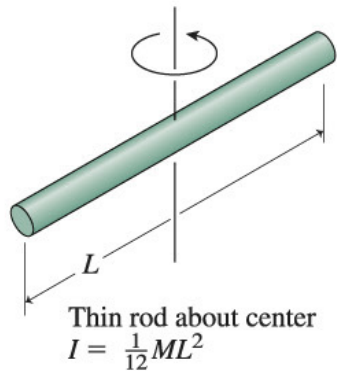


The mass element dm contributes rotational inertia $r^2 dm$.

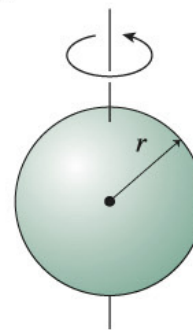


Some Rotational Inertias of Simple Objects

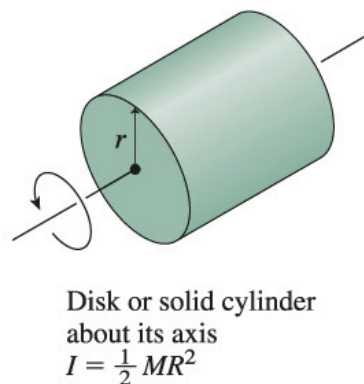
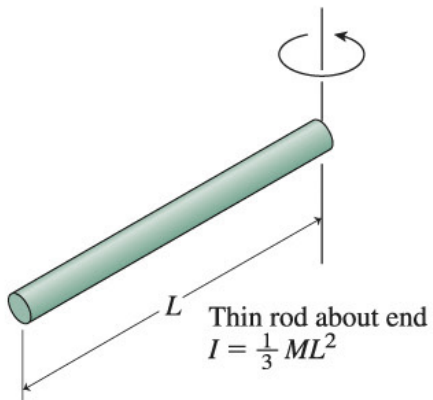
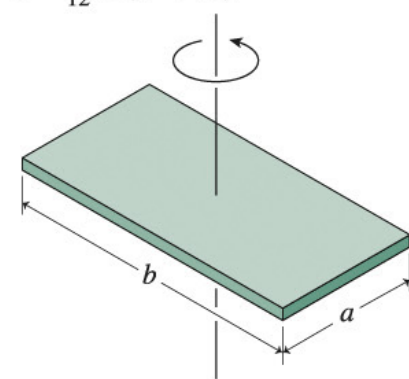
- We really do need to use calculus to figure out rotational inertias of most simple (three-dimensional) geometrical objects



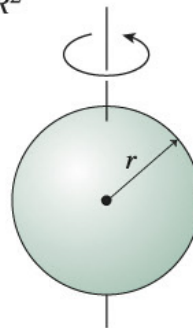
Solid sphere about diameter
 $I = \frac{2}{5} MR^2$



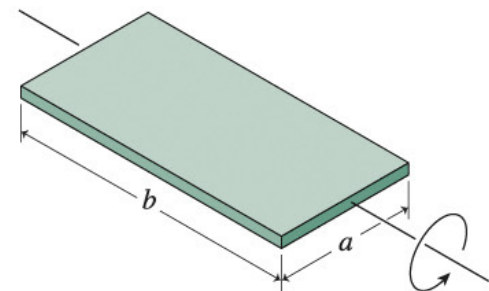
Flat plate about perpendicular axis
 $I = \frac{1}{12} M(a^2 + b^2)$



Hollow spherical shell about diameter
 $I = \frac{2}{3} MR^2$



Flat plate about central axis
 $I = \frac{1}{12} Ma^2$

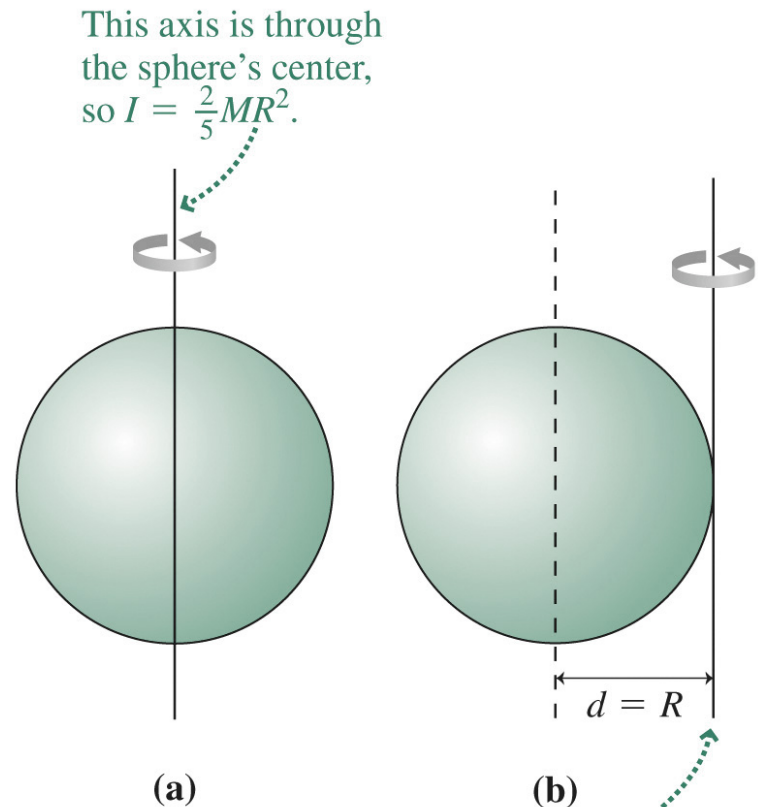


Parallel Axis Theorem

- If we know the rotational inertia I_{cm} about an axis through the center of mass of a body, the **parallel-axis theorem** allows us to calculate the rotational inertia I through any parallel axis.
- The parallel-axis theorem states that

$$I = I_{\text{cm}} + Md^2$$

where d is the distance from the center-of-mass axis to the parallel axis and M is the total mass of the object.



This parallel axis is a distance $d = R$ away from the original axis, so $I = \frac{2}{5}MR^2 + Md^2 = \frac{7}{5}MR^2$.



Example: Hula Hoop Rotational Inertia

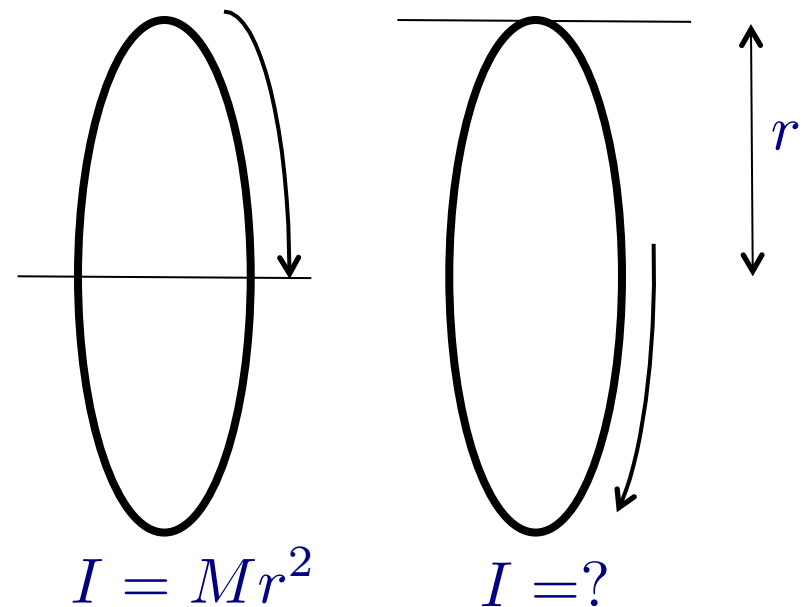
- What is the rotational inertia of a hula hoop of radius r and mass M around its edge?
- Its center of mass is (obviously) at the center of the circle and all of its mass is at the same radius

$$I_{\text{cm}} = Mr^2$$

- The parallel axis theorem gives

$$I_{\text{edge}} = Mr^2 + Mr^2 = 2Mr^2$$

Interesting... It takes twice as much torque to turn a ring around its edge as it takes to turn around its center.

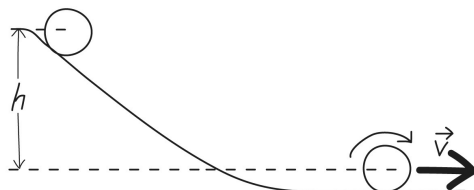
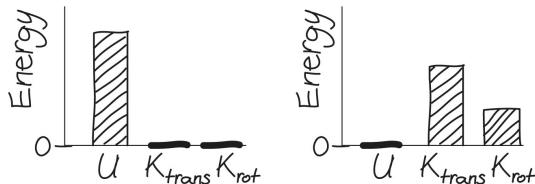


Rotational Energy

- A rotating object has kinetic energy $K_{\text{rot}} = \frac{1}{2} I \omega^2$ associated with its rotational motion alone.
 - It may also have translational kinetic energy: $K_{\text{trans}} = \frac{1}{2} M v^2$.
- In problems involving energy conservation with rotating objects, both forms of kinetic energy must be considered.
 - For rolling objects, the two are related:
 - The relation depends on the rotational inertia.

Example: A solid ball rolls down a hill. How fast is it moving at the bottom?

Energy bar graphs



Equation for energy conservation

$$Mgh = \frac{1}{2} Mv^2 + \frac{1}{2} I \omega^2$$

$$= \frac{1}{2} Mv^2 + \frac{1}{2} \left(\frac{2}{5} MR^2 \right) \left(\frac{v}{R} \right)^2 = \frac{7}{10} Mv^2$$

Solution:

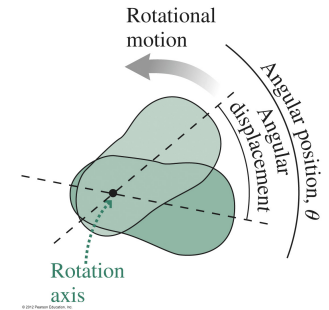
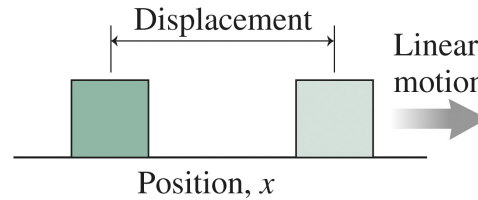
$$v = \sqrt{\frac{10}{7} gh}$$



Summary

- Rotational motion in one dimension is exactly analogous to linear motion in one dimension.

- Linear and angular motion:

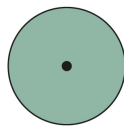


- Analogies between rotational and linear quantities:

Linear Quantity or Equation	Angular Quantity or Equation	Relation Between Linear and Angular Quantities
Position x	Angular position θ	
Speed $v = dx/dt$	Angular speed $\omega = d\theta/dt$	$v = \omega r$
Acceleration a	Angular acceleration α	$a_t = \alpha r$
Mass m	Rotational inertia I	$I = \int r^2 dm$
Force F	Torque τ	$\tau = rF \sin \theta$
Kinetic energy $K_{\text{trans}} = \frac{1}{2}mv^2$	Kinetic energy $K_{\text{rot}} = \frac{1}{2}I\omega^2$	
Newton's second law (constant mass or rotational inertia):		
$F = ma$	$\tau = I\alpha$	

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Rotational inertia, I



Mass closer to axis: lower I



Same mass, farther from axis: greater I

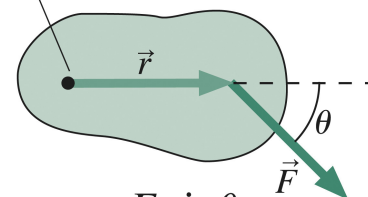
$$I = \sum m_i r_i^2 \rightarrow \int r^2 dm$$

Discrete masses

Continuous matter

Torque, τ

Rotation axis

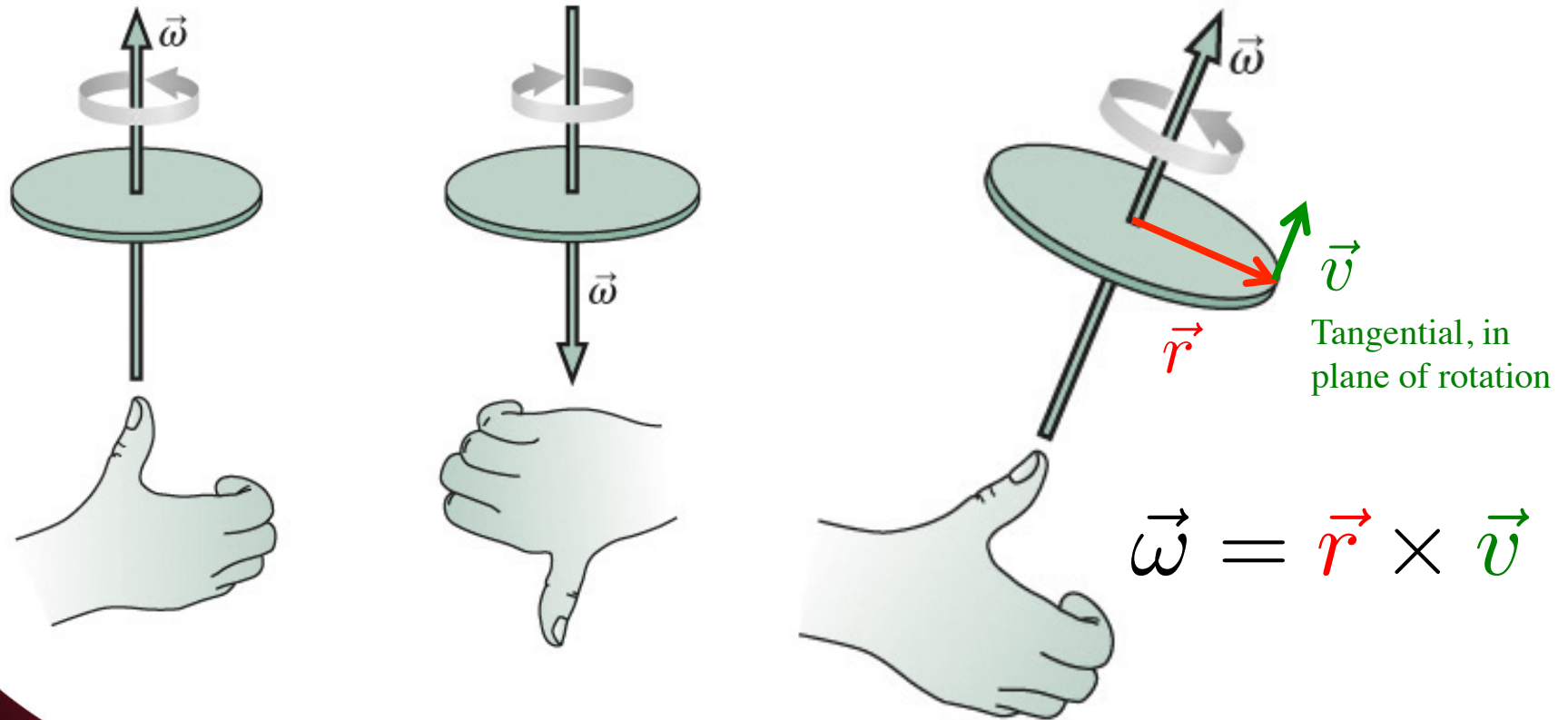


$$\tau = rF \sin \theta$$



Direction of the Angular Velocity Vector

- The direction of angular velocity is given by the **right-hand rule**.
 - Curl the fingers of your right hand in the direction of rotation, and your thumb points in the direction of the angular velocity vector $\vec{\omega}$

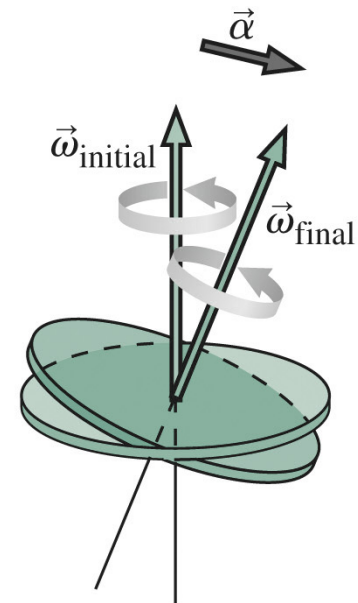
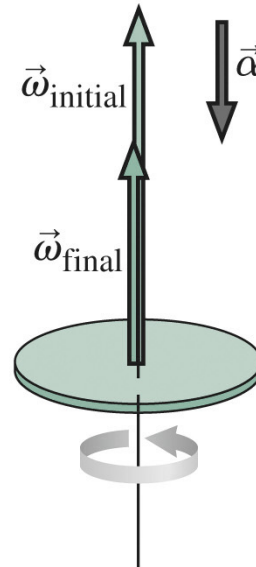
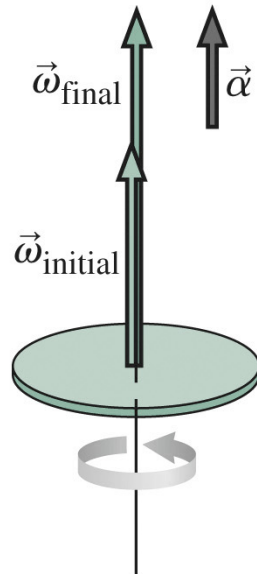


Direction of the Angular Acceleration

- Angular acceleration points in the direction of the change in the angular velocity $\Delta\vec{\omega}$:

$$\vec{\alpha} = \lim_{\Delta t \rightarrow 0} \frac{\Delta\vec{\omega}}{\Delta t} = \frac{d\vec{\omega}}{dt}$$

- $\vec{\alpha}$ can be in the same direction as the angular velocity $\vec{\omega}$, increasing its magnitude but not changing its direction.
- $\vec{\alpha}$ can be opposite the angular velocity, decreasing its magnitude and possibly flipping its sign.
- Or $\vec{\alpha}$ can be in completely different direction than $\vec{\omega}$, changing its direction **and** magnitude.



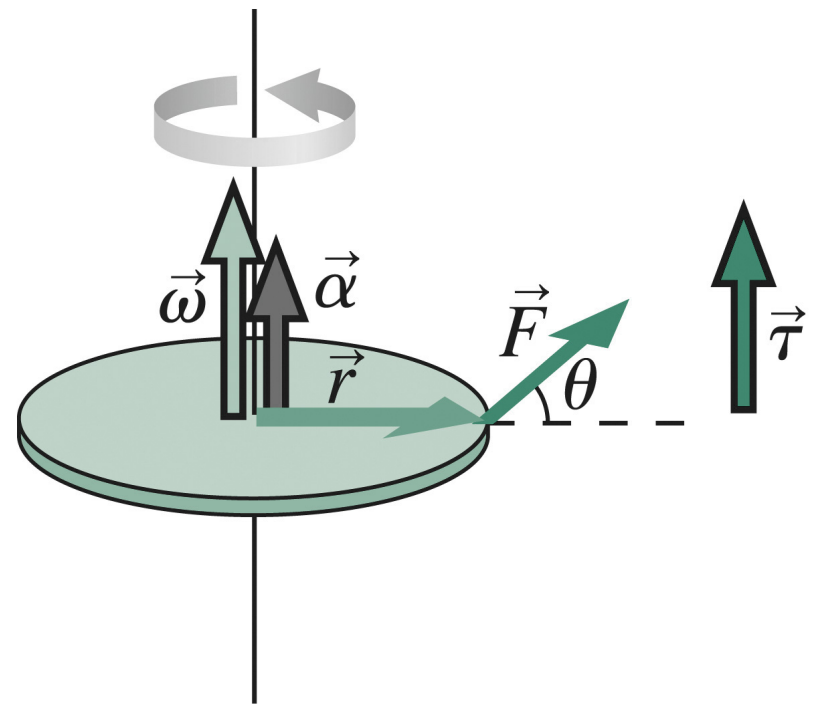
Direction of the Torque Vector

- The torque vector is perpendicular to both the force vector \vec{F} and the displacement vector \vec{r} from the rotation axis to the force application point.
 - The magnitude of the torque is

$$\tau = rF \sin \theta$$

- But torque, radius, and the force are all vectors, and we can more correctly express this as another cross product:

$$\vec{\tau} = \vec{r} \times \vec{F}$$



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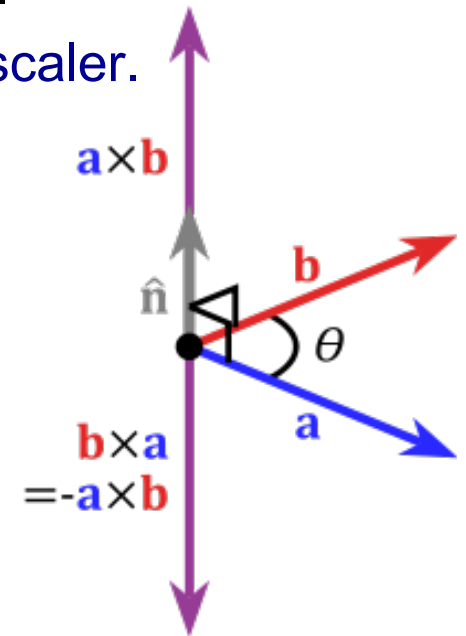
The Cross Product

- The **cross product** or **vector product** is defined as the product of two vectors that produces another vector.
 - This is in contrast to the dot product that produces a scalar.
 - The picture at right shows $\vec{a} \times \vec{b}$ with

$$\vec{b} \times \vec{a} = -\vec{a} \times \vec{b}$$

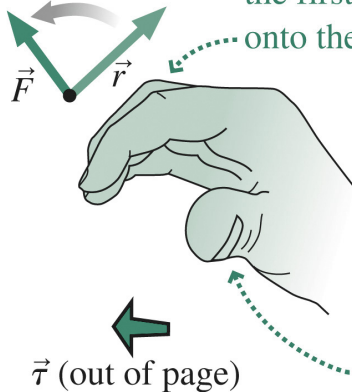
$\vec{a} \times \vec{b}$ is perpendicular to **both** \vec{a} and \vec{b}

- Cross products are **intrinsically** three-dimensional!



Start with the vectors tail to tail.

Curl your fingers in a direction that rotates the first vector (\vec{r}) onto the second (\vec{F}).



$\vec{\tau}$ (out of page)

Then your thumb points in the direction of $\vec{\tau} = \vec{r} \times \vec{F}$.

Some properties of cross products:

$$\vec{A} \times \vec{B} = -\vec{B} \times \vec{A}$$

$$\vec{A} \times (\vec{B} + \vec{C}) = \vec{A} \times \vec{B} + \vec{A} \times \vec{C}$$



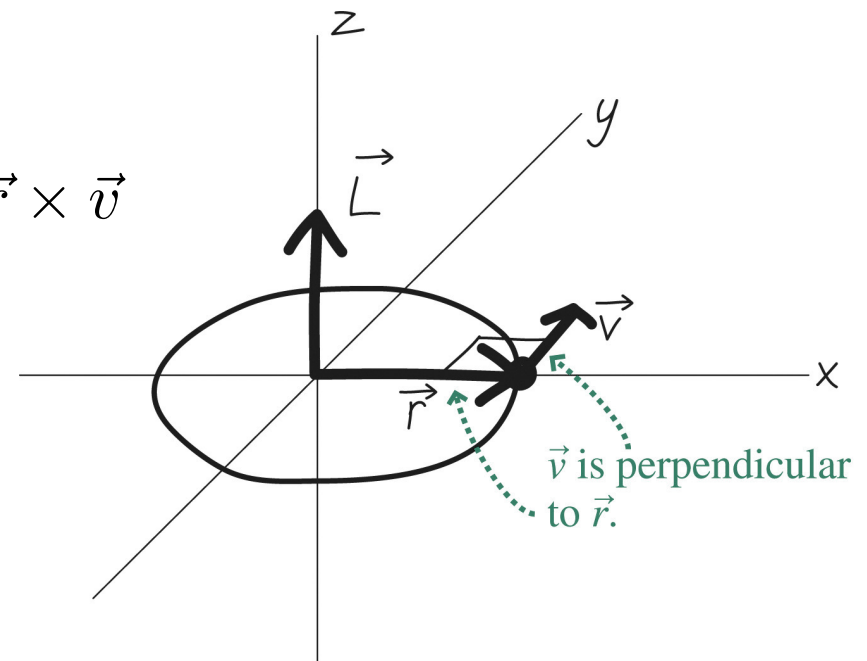
Angular Momentum

- For a single particle, angular momentum \vec{L} is a vector given by the cross product of the displacement vector \vec{r} from the rotation axis with the linear momentum of the particle:

$$\vec{L} = \vec{r} \times \vec{p}$$

- Yes, this looks a lot like $\vec{\omega} = \vec{r} \times \vec{v}$
- For solid objects where we know the rotational inertia, we have another parallel to linear motion:

$$\vec{p} = m\vec{v} \quad \rightarrow \quad \vec{L} = I\vec{\omega}$$



Newton's Law and Angular Momentum

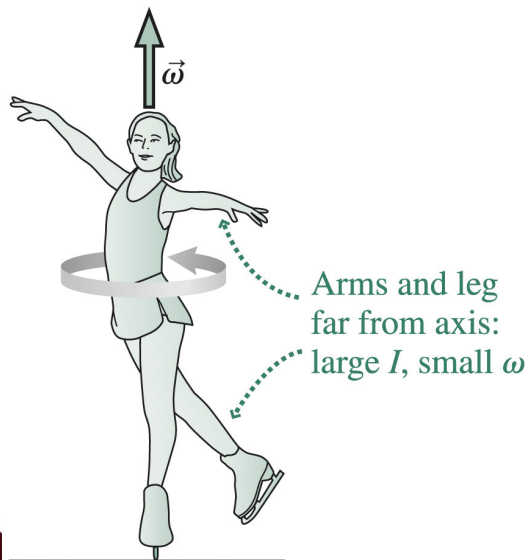
- In terms of angular momentum, the rotational analog of Newton's second law is

$$\vec{\tau} = \frac{d\vec{L}}{dt}$$

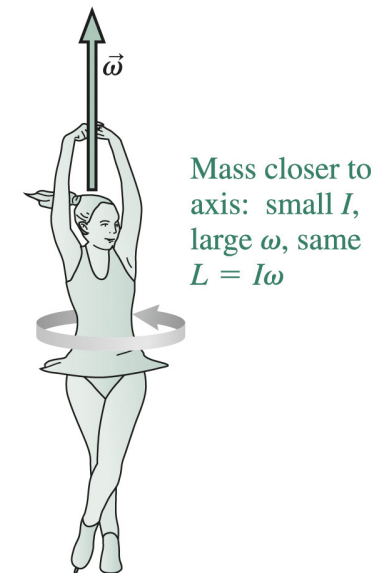
This is very general since rotational inertia I can change!

- A system's angular momentum changes only if there's a non-zero net torque acting on the system.
- If the net torque is zero, then **angular momentum is conserved**.

Changes in rotational inertia I then create changes in angular speed



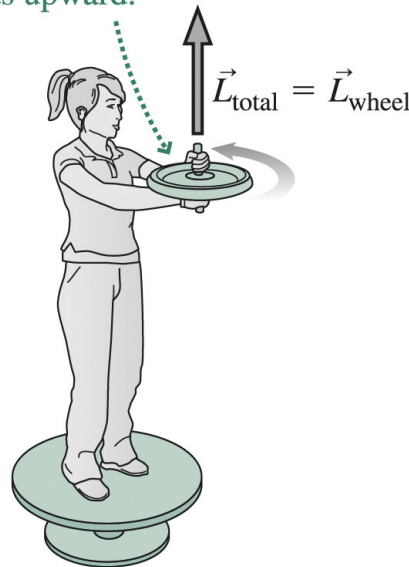
The skater's angular momentum is conserved, so her angular speed increases when she reduces her rotational inertia.



Conservation of Angular Momentum

- The spinning wheel initially contains all the system's angular momentum.
- When the student turns the wheel upside down, she changes the direction of its angular momentum vector.
- Student and turntable rotate the other way to keep the total angular momentum unchanged.

The student stands on a stationary turntable holding a wheel that spins counterclockwise; the wheel's angular momentum points upward.



She flips the spinning wheel, reversing its angular momentum. The total angular momentum is conserved, so turntable and student (ts) must rotate the other way.

