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University Physics 227N/232N Old Dominion University

Conductors, Electric Flux Introduction to Gauss's Law

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Monday, January 27 2014 Happy Birthday to Sam Ting, Rosamund Pike, Patton Oswalt, Lewis Carroll, and Beatrice Tinsley!



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Announcements

- Evening Problem Solving Sessions
 - An opportunity to develop your problem-solving skills!
 - Run by Eric Stacy, estac003@odu.edu, Physics Learning Center, Mondays and Tuesdays 7-9 PM.
 - Similar sessions held 12:30-1:30 on Tuesdays(tomorrow!) in the SCALE-UP classroom (by your valiant TA, Fred Miller).
- Public Talk: The Physics of Football

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- North Cafeteria, Webb Center, Tuesday Jan 28 7 PM (tomorrow!)
- Tim Gay, University of Nebraska, Lincoln
 - Has consulted with NFL Films, ESPN, New York Times, and others
- RSVP to 757-683-3116 or http://www.odu.edu/univevents (TGL14) if you plan to attend.

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Homework Sanity Check

- Have covered material / Today / Wednesday
- Homework Electrostatics II due Thu Jan 31 2014 11:59 PM
 - 20.23: Electric field and force (review)
 - 20.26: Electric field of a proton
 - 20.27: Electric field of two charges, vectors
 - 20.29: Electric field from a wire with constant line density
 - 21.20: Flux

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- 21.23: Flux (sphere)
- 21.25: Gauss's law (basic)
- 21.27: Gauss's law (a little more advanced)
- Tutorial: Field from surface charge on conductor
- There will be a quiz this Friday



Review: Electric Dipole

- An electric dipole consists of two point charges of equal magnitude but opposite signs, held a short distance apart.
 - The dipole is electrically neutral, but the separation of its charges results in an electric field.
 - Many charge distributions, especially molecules, behave like electric dipoles.
 - The product of the charge and separation is the **dipole moment**: *qd*
 - At distances r>>d

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$$E \propto \frac{kqa}{r^3}$$

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(Review: Matter in Electric Fields)

- For a point charge q in an electric field \vec{E} , Newton's second law and the electric force combine to give acceleration: $\vec{a} = q\vec{E}/m$
- A dipole in an electric field experiences a torque that tends to align the dipole moment with the field: $\vec{\tau} = q \vec{d} \times \vec{E}$
- If the field is not uniform, the dipole also experiences a net force.
- The work required to rotate the dipole is $W = -qdE(\cos\theta \cos\theta_0)$ where *q* is the angle between the dipole and the field.
- A dipole in an electric field has a potential energy $U=-q \vec{d}\cdot \vec{E}$



Review: Insulators and Dielectrics

- Materials in which charge isn't free to move are insulators
 - Some insulators contain molecular dipoles, which (as we've seen) experience torques and forces in electric fields.
 - Such materials are called **dielectrics**.
 - Even if molecules aren't intrinsically dipoles, they acquire induced dipole moments as a result of electric forces stretching the molecule.
 - Alignment of molecular dipoles reduces an externally applied field.



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Dielectric Example: Smart Glass

- A modern example of a dielectric application is **smart glass**
 - Glass sandwich with electric dipole polymer filling
 - Adjust transparency by changing external voltage (electric field)



Conductors

- Materials in which charge is free to move are conductors
- In this class, we assume perfect conductors
 - Zero resistance to motion of electrons
 - Many conductors are very good approximations (e.g. Cu, Ag)
 - Superconductors are real examples of perfect conductors
 - We will learn more about them when we learn about magnetism
- Electrons move freely in conductors until electrical forces balance
 - Excess electrons distribute on the surface of a conductor
 - (Excess positive charge "distributes" on the surface of a conductor too)
 - More charge (and higher field) around points





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Conductors and Electric Fields

- Electrons move freely in conductors until electrical forces balance
 - The electric field inside a perfect conductor is always zero
 - Any nonzero field moves electrons until the overall electric field is zero
 - The electric field on the surface of a perfect conductor is always only perpendicular to the surface
 - Any tangential field moves electrons until the tangential field is zero





Onwards to Chapter 21: Gauss's Law

- Represent electric fields using field-line diagrams
- Understand Gauss's law and how it relates to Coulomb's law (charges as "sinks" and "sources")
- Calculate the electric fields for symmetric charge distributions with Gauss's law
- Describe the behavior of charge on conductors in "electrostatic equilibrium"

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Electric Field Lines (or Curves)

- Electric field lines provide a convenient way to generally draw electric fields
 - Each curve's direction at any point is the E field direction
 - Spacing of field lines describes the magnitude of the field
 - Where lines are closer, the field is stronger



Field Lines for Simple Charge Distributions

- There are field lines everywhere, so every charge distribution has infinitely many field lines
- We associate a certain finite number of field lines with a charge of a given magnitude.
- In the diagrams shown,
 8 lines are associated with a charge *q*.
- Field lines of static charge distributions always begin and end on charges, or extend to infinity.

Positive: "sources" Negative: "sinks"

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- How many field lines emerge from any closed surface?
 - Think of + and charges like sources and sinks of electric field lines
 - These are like water (or incompressible) fluid flow
 - Any surface that doesn't enclose a source or a sink
 - must have net flow through it of zero
 - Everything going in also comes out
 - Any surface that encloses only a source
 - must have net flow through it that is outward (positive)
 - Any surface that encloses only a sink

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- must have a net flow through it that is inward (negative)
- Count each field line crossing going outward as +1, each inward crossing as -1.
- We find that the number of field lines crossing any closed surface is proportional to the net charge (sum of charges, or sources and sinks) enclosed.















Activity: Counting Field Lines



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Surface	N(out)	N(in)	N(tot)	Q(enc)
1				
2				
3				
4				

Count these field lines!



Electric Flux

- Electric flux is a precise mathematical version of "number of lines crossing through the surface"
- The electric flux Φ through a flat surface in a uniform electric field depends on the field strength *E*, the surface area *A*, and the angle θ between the field and the normal to the surface.
- Mathematically, the flux over a surface where the electric field intensity is constant is given by

$$\Phi = EA\cos\theta = \vec{E}\cdot\vec{A}$$

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Here \vec{A} is "the" vector whose magnitude is the surface area A and whose orientation is normal (perpendicular) to the surface.



Electric Flux: Flat Surface, Constant E



Electric Flux: Curved Surfaces, Changing E

- When the surface is curved and/or the electric field intensity is not uniform, we can still calculate $\mathbf{flux}\,\Phi$
 - Divide surface into small patches $d\vec{A}$, so small that each patch is (approximately) flat and the field is (approximately) uniform over each

$$d\Phi = \vec{E} \cdot d\vec{A}$$

- Then, as before, we can add up the pieces to get the total flux
 - (The sum becomes an integral)

$$\Phi = \int \vec{E} \cdot d\vec{A}$$

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Although the surface curves and the field varies . . .



Gauss's Law

- We've seen that the electrical flux through a closed surface depends on how much charge is "inside" the surface
 - Makes sense when considering charges as sources and sinks of electrical field (or electrical field lines)
 - Can we write this down mathematically? Yes, the total flux through any surface is proportional to the total charge enclosed $q_{tot,enclosed}$

$$\Phi = \int_{\text{surface}} \vec{E} \cdot d\vec{A} = 4\pi k \ q_{\text{tot,enclosed}} = \frac{q_{\text{tot,enclosed}}}{\epsilon_0}$$

- This is true for any surface and any distribution of charges
- For lots of symmetric problems, we can pick a surface that has the same symmetry, with constant E over the surface area.
- This will make the integral become just a multiplication

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Gauss's Law: Example (Point Charge)

- The electric field from a point charge points radially outward
- A sphere of radius r centered on the charge has a surface that's always perpendicular to the electric field



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$$\vec{E} \cdot d\vec{A} = E \, dA \, \cos \theta$$

 $\vec{E} \perp d\vec{A} \Rightarrow \cos \theta = 1$
 $\vec{E} \cdot d\vec{A} = E \, dA \, \text{over}$

entire surface of sphere

$$\Phi = \int_{\text{surface}} \vec{E} \cdot d\vec{A} = 4\pi k \ q_{\text{tot,enclosed}} = \frac{q_{\text{tot,enclosed}}}{\epsilon_0}$$

$$\Phi = E A = E \left(4\pi r^2\right) = 4\pi kq \qquad E = \frac{\kappa q}{r^2}$$

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